An introduction to Statistical Energy Analysis



BREKKE STRAND

Guest lecture 11/12/19

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Myself

- 2002-2005 BSc in Telecom Eng, Politecnico Milano
 - voice activity detector GSM networks, Siemens
- 2005-2008 MSc in Sound Eng and Design, Politecnico
- 2006-2008 MSc in Sound and Vibration, Chalmers
 - parametric array loudspeakers
- 2008-2013 PhD in Technical Acoustics, KTH
 - SEA, porous materials, variational principles

Myself

• 2013-2016, Consultant, Lloyd's Register Consulting

- N&V; marine sector
- 2016-2017, Consultant, Creo Dynamics
 - N&V; R&D, ultrasonics, nonlinear acoustics, Volvo cars.
- 2017-2019, Consultant, WSP
 - N&V; urban and industrial sectors.
- 2019 today, Consultant, Brekke & Strand (Malmö)

BREKKE STRAND

This lecture

- Material produced at KTH
 - internal courses
 - seminars/lectures within Mid-Mod EU project
 - together with my supervisor Svante Finnveden (former professor at KTH, now at ACAD)



• Rocket to the moon (Smith, 1962)

- some 200k modes
- modelling effort



- low frequency: measurements, FEM
- high frequency: today's topic
- Definition of high frequency

- Broadband excitation of a structure
- Not too high damping \Rightarrow reverberant field
- Peaks, natural frequencies ⇔ modes



Fahy, Gardonio



- Clear peaks
- Peaks characteristic of the structure

- Not clear peaks
- Peaks dependent on small geometric variations, discontinuities



 Magnitudes of the structure-borne frequency response functions of 98 nominally identically cars (Isuzu Rodeo); force applied on the wheel hub, microphone at the driver's seat. (Kompella, Bernhard)



12 measurements on the same car (Kompella, Bernhard)



Vibroacoustic responses of 141 identical beer cans (Fahy)



Figure 2.12 Standard deviation of ten measurements of airborne transmission through a concrete floor and ten measurements of structural transmission between a wall and a floor. ———, airborne transmission; — — —, structure-borne transmission (Craik and Steel, 1989; Craik and Evans, 1989)

Vibroacoustic responses of *identical* building elements in same building

Motivation - What did we see?

- What is changing between various cars from the same production line?
- What is changing between measurements on the same car?

- We may have a feeling about what high frequency means
- We may have a measure for this



Input power to a plate with nominal dimensions (left), to 8 plates with nominal dimensional + std(0.02) on h (right)





Motivation - What can we say?

- Two problems with high frequency models:
 - Technical: small mesh compared to wavelength (FE)
 - profound: chaotic response
 - deterministic;
 - apparently random;
 - unpredictable.

Motivation - What can we say?

- Useful approach i.e. not the only but a useful one for specific needs:
 - quick and efficient
 - accounting for statistical variation
 - employing a measure resistant to variations

Answer

- Statistical Energy Analysis
 - statistical
 - energy
 - analysis
 - Seems Easy, Ain't
- Frequency, space, ensemble average response

Answer





Input power to 8 plates with nominal dimension + std(0.02) (black); SEA prediction (red)

Formulation



nowadays models may look like CAD drawings...

Formulation: 0th step



- To start
- a structure
- identify sources
- linear motion

Formulation: 1st Step

- Subdivide structure/vibroacoustic field into SEA Elements
 - elements may be substructures
 - elements may be elements of the response
 - elements are quite large
 - very difficult task!

Roof bending
Engine

Roof longitudinal
Cabin air volume

Formulation: 2nd step

Power balance for element i

$$P_{in}^{i} = P_{diss}^{i} + \sum_{j \neq i} P_{coup}^{ij}$$

- SEA built upon conservation of energy!
- Now we need to define the terms in the summation

Formulation: 3rd step

Dissipated power

$$P_{diss}^i = \eta \omega E$$

- η damping loss factor
- ω angular frequency
- E total vibroacoustic energy
- nothing really new (Lord Rayleigh)

Formulation: 4th step



$$P_{coup}^{ij} \propto \left(\frac{E_i}{\Delta N_i} - \frac{E_j}{\Delta N_j}\right)$$

- ΔN_i number of modes in a band
- introduce:

 $n_{i} = \frac{\Delta N_{i}}{\Delta \omega}$ modal density [1/rad/s] $\hat{e}_{i} = \frac{E_{i}}{n_{i}}$ modal energy (power!) [W]

-
$$C^{ij} = \omega n_i \eta_{ij} = \omega n_j \eta_{ji} = C^{ji}$$
 conductivity [-]

$$P_{coup}^{ij} = C^{ij} \left(\hat{e}_i - \hat{e}_i \right)$$

Formulation: 5th Step

- Modal power potentials $\hat{e}_i = E_i / n_i$ [W]Dissipated Power $P_{diss}^i = M_i \hat{e}_i$ [W]
- Modal overlap factor $M_i = \eta_i \omega n_i$ [-]
- Coupling power $P_{coup}^{1,2} = C(\hat{e}_1 \hat{e}_2)$ [W]

C – Conductivity [-]; n – modal density [1/(rad/s)]; η - loss factor [-].

Power balance: $P_{dis} + P_{coup} = P_{in}$:

$$\begin{bmatrix} M_1 + \sum_j C^{1j} & -C^{12} & -C^{13} & \cdots \\ -C^{12} & M_2 + \sum_j C^{2j} & -C^{23} & \cdots \\ \vdots & \vdots & \ddots & \vdots \\ & & & \ddots \end{bmatrix} \begin{bmatrix} \hat{e}_1 \\ \hat{e}_2 \\ \vdots \end{bmatrix} = \begin{bmatrix} P_{in}^1 \\ P_{in}^2 \\ \vdots \end{bmatrix}$$
 Linear system of equations

SEA quantities: Energy

- Energy as dependent variable, unknown
- Frequency, space, ensemble averaged energies (ergodic assumption)
- SEA deals with:
 - resonant motion
 - reverberant wave motion
- Strain energy (density) = Kinetic energy (density)
- Energy density <u>constant</u> across the element

SEA quantities: Energy



Retrieve vibroacoustic quantities!

 Modal density: *expected* number of resonance per unit frequency

$$n_i = \frac{\Delta N_i}{\Delta \omega}; \qquad n = \frac{1}{\delta \omega}$$



expected difference
 between natural frequency
 [rad/s]



Real part of mobility for statistic ie diffuse conditions:

 $\operatorname{Re}\left\{Y\right\} = \frac{\pi}{2M} \frac{\Delta N}{\Delta \omega}$

Figure 5.63 Measured mobility of a $6 \times 4 \text{ m}^2 \times 200 \text{ mm}$ thick concrete floor: real part; … imaginary part; – – – theoretical infinite plate mobility (real); + predicted modal resonance frequencies.

Craik, Sound transmission through buildings using statistical energy analysis Fahy, Noise and vibration

- Statistics implies little/no difference between various boundary conditions and/or geometries!
- Infinite structures just as good for some applications.



Figure 5.63 Measured mobility of a $6 \times 4 \text{ m}^2 \times 200 \text{ mm}$ thick concrete floor: real part; … imaginary part; – - – theoretical infinite plate mobility (real); + predicted modal resonance frequencies.

Table 5.3. Number of modes and modal densities

Rod (lon- gitudinal)	$N = k_L l/\pi = \omega l/(c_L \pi)$	$\Delta N \big/ \Delta \omega = l \big/ \big(c_L \pi \big)$	
Beam (bending)	$N = k_B l / \pi = \sqrt{\omega} l / \left(1.7 \sqrt{c_L h} \right)$	$\Delta N / \Delta \omega = k_B l / (2\pi \omega)$ $= l / (3.4 \sqrt{\omega c_L h})$	
Plate (bending)	$N = k_B^2 S / 4\pi = \omega S / (3.6c_L h)$	$\Delta N / \Delta \omega = k_B^2 S / (4\pi \omega)$ $= S / (3.6c_L h)$	
Acous- tic volume	$N = k_0^3 V / 6\pi^2 = \omega^3 V / (6\pi^2 c_0^3)$	$\Delta N / \Delta \omega = k_0^2 V / (2\pi^2 c_0)$ $= \omega^2 V / (2\pi^2 c_0^3)^{-1}$	
Ring (radial)	$N = 2k_B a = 3.7\sqrt{\omega} a / \sqrt{c_L h}$	$\Delta N / \Delta \omega = k_B a / \omega$ $= 1.9 a / \sqrt{\omega c_L h}$	
Thin- walled cy- lindrical shell, v < 1	$N\approx 3\sqrt{3}N^{3/2}/(2\pi h)$	$\Delta N/\Delta \omega \approx 2\sqrt{\omega} la^{3/2} / (1.6hc_L^{3/2})$	
Thin- walled cy- lindrical shell, v > 1	$N \approx \sqrt{3} la \omega / (c_L h)$	$\Delta N / \Delta \omega \approx \sqrt{3} la / (c_t h)$	

 k_B, c_B : wavenumber and sound speed for bending wave,

l length, *h* thickness, *S* area, *V* volume, *a* radius and $v = \omega a/c_L$.

Cremer, Heckl, Pettersson, Structure-borne sound

Modal overlap factor

$$\begin{split} P^{i}_{diss} &= \eta \omega E = M \hat{e} \\ M &= \eta \omega n; \qquad M = \frac{\eta \omega}{\delta \omega} \end{split}$$

- quantifies dissipation
- quantifies probability of a resonance

Harmonic forced response of MDOF and continuous systems



Many modes, many natural frequencies, many resonances.

Response dominated by resonant contributions.

Half power bandwidths generally increase with increasing frequency: resonance peaks overlap.

 $M = \frac{\eta\omega}{\delta\omega}$ $M = \eta \omega n;$

(Brian Mace)



Fig 2-4. Real part of point mobility for a car tyre. (The real part of the mobility multiplied by the squared force magnitude gives the power injected into the structure by a point force.)

Finnveden, Energy methods (lecture notes KTH course)

approx. dampig loss factor = 7 Hz / 140 Hz = 0,05

 Table 9.1 Comparison of Commonly Used Damping Measures

Dissipation Descriptor	Symbol	SI Units	Relation to η
Loss Factor	η	-	η
Quality Factor	Q	-	$1/\eta$
Critical Damping Ratio	ζ	-	$1/\eta$
Reverberation Time	T_r	sec.	$2.2/f\eta$
Decay Rate	DR	dB/sec.	$27.3 f\eta$
Logarithmic Decrement	δ	nepers/s	$\pi f\eta$
Wave Attenuation*	γ	nepers/m	$\pi f\eta/c_g$
Mechanical Resistance**	R	N-s/m	$2\pi f\eta M$
Damping Bandwidth (Half-Power)	BW	Hz	$f\eta$
Imaginary Part of Modulus $\mathbf{E}_r + j\mathbf{E}_i$	\mathbf{E}_{i}	N/m ²	$\mathbf{E}_r \eta$
Acoustical Absorption Coefficient***	α	-	$(8\pi fV/cA)\eta$

* c_g is group velocity for system in m/s

- ** M_g is the system mass
- ******* A is the area of walls, V is room volume, c is speed of sound

SEA quantities: Conductivity

$$P_{coup}^{ij} = C^{ij} \left(\hat{e}_i - \hat{e}_i \right)$$

Potential Flow -Thermal analogy



Thermal system: heat and temperature. SEA system: power and vibrational energy.

Temperature ~ energy per mode. Total thermal energy ~ total vibrational energy. Heat flows from "hot" to "cold". The bigger the temperature difference, the bigger the heat flow.

(Brian Mace)
$$P_{coup}^{ij} = C^{ij} \left(\hat{e}_i - \hat{e}_i \right)$$

- A statistical approach based on a low population!
 - Thermodynamics: Avogadro's number, 10^23
 - SEA: vibroacoustic natural modes, 0 10^4?

$$P_{coup}^{ij} = C^{ij} \left(\hat{e}_i - \hat{e}_i \right)$$



(Brian Mace)

when calculating C

- Compact support
 - only direct couplings
- Travelling wave estimate:
 - no transfer of energy back
 - regardless of rest of the structure







- Where SEA should work due to chaotic response (M>=1), then it predicts well the conductivity (coupling loss factor).
- i.e. extending the receiving structure to infinity is sensible!



pin



$$\begin{array}{ll} 1. \quad E = 2\langle e_p \rangle V = 2\frac{1}{2}\frac{\tilde{p}^2}{\rho_0 c^2}V & [J]; & \text{speed of sound c} \\ \text{speed of sound c} \\ \text{air density rho_0} \end{array}$$

$$\begin{array}{ll} 2. \quad \hat{e} = \frac{E}{n} & \left[\frac{J}{s}\right] = [W]; \\ 3. \quad n = \frac{k^2 V}{2\pi^2 c} & [s]; & \text{wavenumber } \mathbf{k} = \omega/c \\ 4. \quad M = \omega\eta n = \frac{k^2 A}{8\pi^2} & [-]; & A = \alpha S; & \text{partition areas S} \\ 5. \quad C = \left(\frac{W_t}{\hat{e}_1}\right)_{\substack{C^{12} \\ M_2 \to 0}} = \frac{\tau W_i}{\hat{e}_1} & [-]; & \begin{array}{c} \text{incident power W_i} \\ \text{transmitted power W_i} \\ \text{transmitted power W_i} \\ \text{transmitted power W_i} \\ \text{transmission factor tau} \end{array}$$

$$\begin{array}{c} 6. \quad \tau = \frac{W_t}{W_i}; & R = 10\log\frac{1}{\tau}; & W_i = \frac{k^2 S}{8\pi^2}\hat{e}_1; \\ 7. \Rightarrow C = \tau \frac{k^2 S}{8\pi^2} \end{array}$$

Room acoustics exercise

$$\begin{bmatrix} M_{1} + C & -C \\ -C & M_{2} + C \end{bmatrix} \begin{bmatrix} \hat{e}_{1} \\ \hat{e}_{2} \end{bmatrix} = \begin{bmatrix} P_{in} \\ 0 \end{bmatrix};$$

$$\hat{e} = \frac{eV}{n} = \frac{2\pi^{2} \langle \tilde{p}^{2} \rangle}{\omega^{2} \rho_{0} c};$$
Standard texts
1. $L_{p1} = L_{w} + 10 \log \left(\frac{\Gamma}{4\pi r^{2}} + \frac{4}{A'} \right);$ direct field is not SEA
2. $L_{p2} = L_{p1} - 10 \log \left(\frac{A}{\tau S} + 1 \right)$ typical approximation

- Symmetric formulation
- Sub-structuring
- Distinction between coupling and damping losses
- (nothing really new but a good framework!)
 - (good news: you already know a bit of SEA!)







• Coincidence frequency of a plate





Leppington, The acoustic radiation efficiency of rectangular panels, Proceedings of the Royal Society of London A, 382 (1982)

 Radiation efficiency (k_f, fluid wavenumber; k_s structure wavenumber):

Assuming fc >> 1st eigenfrequency of plate

$$\sigma(v) = \begin{cases} \frac{c(a+b)}{2\pi^2 v k_f a b (v^2 - 1)^{\frac{1}{2}}} \left[\ln\left(\frac{v+1}{v-1}\right) + \frac{2v}{v^2 - 1} \right], v > 1\\ \left(0.5 - 0.15 \frac{a}{b} \right) \sqrt{k_f a}, v = 1\\ \left(1 - v^2 \right)^{-\frac{1}{2}}, v < 1 \end{cases}$$

where *a* and *b* are the side lengths, so that S = ab, $a \le b$, and

$$\nu = k_{s}/k_{f}.$$

Leppington, The acoustic radiation efficiency of rectangular panels, Proceedings of the Royal Society of London A, 382 (1982)

- Mass-law, forced (non-resonant) motion below coincidence frequency (this is NOT SEA!):
 - insert this *indirect coupling* between two rooms as direct coupling!
 - (what would be an example of actual direct coupling between two rooms?)



- Mass-law, forced (non-resonant) motion below coincidence frequency (this is NOT SEA!):
- insert indirect coupling between two rooms as direct coupling
- Transmitted and incident power in a room, diffuse field

$$W_{\text{trans}} - W_{\text{inc}}\tau - \frac{Ec_0S}{4V}\tau.$$

Coupling loss factor (as per definition):

$$\eta_{12} = \frac{c_0 S \tau_{12}}{8 \pi f V_1}.$$

Craik, Sound transmission through buildings using statistical energy analysis



non-resonant transmission coefficient by Leppington

Figur 43 Principiellt utseende för reduktionstal hos en enkelvägg

There are many approximations to this equation which are widely used (and misused) in building acoustics. The normal incidence sound reduction index, R_0 , can be given by (Beranek and Ver, 1992)

$$R_{0} - 10\log\left[1 + \left(\frac{2\pi f\rho_{s}}{2\rho_{0}c_{0}}\right)^{2}\right] \simeq 20\log\frac{2\pi f\rho_{s}}{2\rho_{0}c_{0}} - 20\log f\rho_{s} - 42$$
(4.23)

and can be obtained from consideration of forced waves at normal incidence on an infinite plate.

Variations of this are the field incidence mass law,

$$R_{f_{1}} - R_{0} - 5 - 20 \log f \rho_{s} - 47 \qquad (4.24)$$

Craik, Sound transmission through buildings using statistical energy analysis

 Use eq. 4.24 to obtain R_f (non-resonant sound reduction index) and use it to compute tau₁₂ and then eta₁₂ i.e. the coupling loss factor describing the non-resonant transmission between two acoustic volumes (two slides ago).

- Finalise: get sound pressure levels in the two rooms.
 - Functions of modal energy, retrieve pressure.
- Check formula in e.g. ISO standard 16283-1:2014 to retrieve R from sound pressure levels in the two rooms, wall area and absorption area of the receiving room.

• Done!

Craik, Sound transmission through buildings using statistical energy analysis

Summary

- Helmholtz number, He = kl = $2\pi/\lambda$ l (wavenumber k, typical dimension l).
- Standing wave at least, $I = \lambda/2$. Hence, $He = \pi$.

Classification	Frequency region	Character
No-modes region	He << π	No wave propagation. Discrete mechanical system.
Low-modes region	He $\approx \pi$	A few eigenmodes dominate. A complete mathematical description is possible (analytical / FEM).
Many-modes region	He >> π	A large number of eigenmodes control the vibroacoustic response. Analyses is only practicable / sensible using energy methods

 The energy method presented today is SEA, which has its range of validity – i.e. strengths and weaknesses.

Validity of SEA

Careful! If losses are too high, then energy decay will occur! Perhaps $M \approx 1$ works better.

- N>>1; many modes in a given frequency band
- M>>1; we need disorder, i.e. overlapping modes
- low energy decay within an element == constant energy density within an element
- weak coupling: Cij/Mi << 1, i.e. flow of exchanged vibroacoustic energy small compared to internal losses. Think how Cij is estimated.

Why SEA?

- Only method for high frequency vibroacoustics of complex systems supported by commercial software
- Efficient (numerically cheap)
 - less demand on input data
 - immediate answers
 - impossible problems
- We are ignorant of many things
 - why bother with unrealistic precision?
- mean response of "typical" product
- Alternate view point



Figure 2.12 Standard deviation of ten measurements of airborne transmission through a concrete floor and ten measurements of structural transmission between a wall and a floor. ———, airborne transmission; — — —, structure-borne transmission (Craik and Steel, 1989; Craik and Evans, 1989)

Why SEA?

- Successful with:
 - satellites, launch vehicles
 - airplane
 - ships
 - buildings
 - vehicles

Why not SEA?

- No details in results
 - poor frequency resolution
 - no spatial resolution
- Software
 - lack of routines for parameter estimation (not so true actually)
 - lack of bench mark examples
- Connection to deterministic methods
- Not universally applicable (narrowband excitation)
 - how accurate?
 - valid for this problem?

Why not SEA?

- In none of my jobs as consultant I have used SEA!
- Instead:
- Semi-empirical method based on measured transfer functions
- Bastian stripped down first-level SEA based on reduction index and impact sound isolation
- Just guess high frequency performance based upon experience.

Why not SEA?



- As seen today, SEA is often a point of view, a framework and in that sense I have carried it with me all the time.
- Plus, room acoustics builds on SEAs thinking, common software (Bastian, EN 12354) too.

References

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